Analysis of strange-mode instability with time-dependent convection in hot massive stars

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Abstract. We carry out nonadiabatic analysis of strange-modes in hot massive stars with timedependent convection (TDC) for the first time. Although convective luminosity in envelopes of hot massive stars is not as dominative as in stars near the red edge of the classical Cepheid instability strip in the Hertzsprung-Russell (H-R) diagram, we have found that the strange-mode instability can be affected by the treatment of convection. However, existence of the instability around and over the Humphreys-Davidson (H-D) limit is independent of the treatment. This implies that the strange-mode instability could be responsible for the lack of observed stars over the H-D limit regardless of uncertainties on convection theories.

Keywords. stars: Hertzsprung-Russell diagram, stars: oscillations (including pulsations)

"Strange" modes with extremely high growth rates appearing in very luminous stars with $L/M \gtrsim 10^4 L_{\odot}/M_{\odot}$ have significantly different characteristics from ordinary pand g-modes (Wood 1976, Shibahashi & Osaki 1981). Stability of strange-modes in hot massive stars has been analyzed with frozen-in convection (FC) by Glatzel & Mehren (1996), Godart *et al.* (2011), Saio *et al.* (2013), and others. In envelopes of hot massive stars, convective luminosity is not as dominant as in stars near the red edge of the classical Cepheid instability strip. But the strange-modes are excited at the convection zones, and we cannot definitely conclude that effects of convection are negligible.

We construct stellar models with X = 0.70, Z = 0.02 by using MESA (Paxton *et al.* 2011), and analyze their radial modes with the nonadiabatic code developed by Sonoi & Shibahashi (2012). We use the time-dependent convection (TDC) formulation derived by Unno (1967) and developed later by Gabriel *et al.* (1974) and by Gabriel (1996). This TDC theory has already been independently implemented by Grigahcène *et al.* (2005). Figure 1 shows the results for $60 M_{\odot}$ with two types of FC; zero Lagrangian and Eulerian perturbations of convective luminosity, $\delta L_{\rm C} = 0$ and $L'_{\rm C} = 0$. The ascending sequences such as A1 and A2 correspond to ordinary modes, while the descending ones such as D1, D2 and D3 are strange-modes. On the whole, instability is suppressed in the $L'_{\rm C} = 0$ case compared with the $\delta L_{\rm C} = 0$ case. The pulsations with $L'_{\rm C} = 0$ are damped due to the substantial gradient of convective luminosity near the boundaries of convective layers, while there is no convective damping with $\delta L_{\rm C} = 0$.

Figure 2 shows the result with TDC. The degree of the instability is in between the two types of FC shown in Fig. 1. While unstable modes of D1 and D2 are excited around the convective layer caused by the Fe opacity bump, where the ratio of the convective to the total luminosity $L_{\rm C}/L_{\rm r}$ is ~ 10⁻¹, the sequence of D3 is pulsationally unstable invariably in all the three treatments, as seen in Figs. 1 and 2. These unstable modes of D3 are excited around the layer of He opacity bump, where $L_{\rm C}/L_{\rm r}$ is negligibly small, and hence we conclude that those modes are definitely unstable regardless of uncertainties of convection theories. We note that the instability of D3 appears around and above



Figure 1. Modal diagrams of radial modes for $60 M_{\odot}$ with FC of zero Lagrangian (*left panel*) and Eulerian (*right panel*) perturbations of convective luminosity. The open and filled circles denote stable and unstable modes, respectively, and the vertical dashed line in each panel shows the cross point of the evolutionary track and the H-D limit.



Figure 2. Results with TDC. Left: Same type of modal diagram as Fig. 1. Right: Instability ranges on the evolutionary tracks derived with TDC on the H-R diagram. The thin dashed lines are evolutionary tracks, and the solid parts correspond to evolutionary stages having at least one unstable mode. Instability of D3 is particularly indicated with the thick solid lines.

the H-D (Humphreys & Davidson 1979) limit. The resultant phenomena by the strangemode instability have been proposed to be pulsationally driven mass loss by nonlinear analyses (e.g., Grott *et al.* 2005) and by observation (Aerts *et al.* 2010). More detailed investigations are worth doing.

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