Distribution of basal melting and freezing beneath tributaries of Ice Stream C: implication for the Holocene decay of the West Antarctic ice sheet

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ABSTRACT. Ice-stream tributaries connect the relatively slow-moving interior of the West Antarctic ice sheet (WAIS) with the fast-flowing Siple Coast ice streams. Basal water underneath these ice streams reduces basal resistance and enables the fast motion of the ice. Basal melting being the only source for this water, it is important to include the distribution of basal melting and freezing into numerical models assessing the stability of the WAIS. However, it is very difficult to constrain its distribution from existing field observations. Past borehole observations confirmed the presence of a wet bed at Byrd Station in the WAIS interior and at different locations within Siple Coast ice streams. However, the recent discovery of a 12-25 m thick sediment-laden bubble-free basal ice layer at the UpC boreholes indicates that basal freezing is either currently occurring or had occurred upstream during the last glacial-interglacial cycle. We use a flowline model of ice thermodynamics to assess and quantify the spatial and temporal distribution of basal melting and freezing beneath Ice Stream C tributaries, taking into account the geothermal flux, shear heating and heat conduction away from the bed. Under the assumption that the ice was moving over a weak bed ($\tau_{\rm b} = 1-10$ kPa) our model is able to reproduce a layer of frozenon ice similar in thickness to the UpC "sticky spot" basal ice layer. Increased basal melting in the early Holocene possibly could have initiated the Holocene decay of the WAIS, whereas increased freezing rates over the past few thousand years could have decreased the amount of basal water in the system, resulting in a strengthening of the bed. This is consistent with current force-budget calculations for ice-stream tributaries and with observed stoppages and slow-downs of ice streams.

INTRODUCTION

The possible future instability of the West Antarctic ice sheet (WAIS) and its effects on global sea-level rise have been a major focus of Antarctic research for the last several decades. Field and remotely sensed data revealed a complex structure in the WAIS drainage system (Joughin and others, 1999) and enabled us to better elucidate the basal processes that permit fast ice-stream motion under low driving stresses (Alley and others, 1987a, b, 1989; Blankenship and others, 1987; Engelhardt and others, 1990; Engelhardt and Kamb, 1997; Anandakrishnan and others, 1998; Bell and others, 1998; Tulaczyk and others, 1998, 2000a, b; Kamb, 2001). High basal water pressures and a layer of weak till with a large porosity seem to play the key role in facilitating the fast ice streaming in West Antarctica (Engelhardt and others, 1990; Harrison and others, 1993, 1998; Echelmeyer and others, 1994; Engelhardt and Kamb, 1997; Tulaczyk and others, 1998, 2000a, b; Kamb, 2001). It is therefore clear that the spatial and temporal availability of basal water has to be incorporated into models simulating the present and future behavior of the WAIS (see, e.g., Budd and others, 1984, 1998; MacAyeal, 1993; Fastook and Prentice, 1994; Payne, 1998, 1999; Hulbe and MacAyeal, 1999; Hulbe and others, 2000; Raymond, 2000).

Borehole observations in the interior of the WAIS (Robin, 1983) and in the Siple Coast ice streams (Engelhardt and others, 1990) revealed a wet bed with ice at, or close to, its pressure-melting temperature. Recent drilling at the UpC "sticky spot" revealed a 12-25 m thick sediment-laden, bubble-free layer of basal ice (personal communication from H. Engelhardt, California Institute of Technology (Caltech) field season 2000/01), henceforth referred to as the basal ice layer. The appearance and composition of the layer are consistent with its origin by basal freeze-on (Vogel and others, unpublished stable-isotope data). The presence of this basal ice layer indicates that basal freeze-on, rather than melting, prevailed for at least several thousand years, upstream of the UpC area. Lacking direct observations from ice-stream tributaries, which channel ice toward the UpC area, we used a thermodynamical numerical model to assess their basal energy balance. Our goal is to test whether accretion of the thick basal ice layers could have occurred as the ice, found now in the UpC area, traveled through the tributaries of Ice Stream C.

ICE STREAM C CATCHMENT AREAS

Ice Stream C (Fig. 1) is located in the central part of the Siple Coast. The Ice Stream C trunk is separated from its southern neighbor Whillans Ice Stream (formerly Ice

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Fig. 2. Profiles of the four simulated flowlines converging toward the UpC area of Ice Stream C: (a) CN-NW; (b) CN-S; (c) CS-N; and (d) CS-S. The upper diagram of each flowline shows surface and bed elevation, while the lower diagram shows the ice velocity used in the model and the observed ice velocity (dotted where it differs from the model velocity). In the upper part of the ice stream we have arbitrarily adjusted ice velocities to account for the faster ice flow expected prior to the shut-down of Ice Stream C.

Stream B) by the interstream ridge B/C, and from its northern neighbor Ice Stream D by Siple Dome. Having two main tributaries, Ice Stream C draws its ice from the northeastern WAIS divide common with the Amundsen Pine Island Glacier catchment area and from the southeastern WAIS divide common with the Weddell Sea catchment area. Originating in the Byrd Subglacial Basin (BSB; up to 2500 m below sea level), the northern tributary draws its ice from the more northeasterly ice divide. Converging from the surrounding mountain blocks into the BSB, the ice reaches a maximum thickness of > 4000 m, before it speeds up and emerges following the Bentley Subglacial Trench

Fig. 1. Bed topography (a), surface elevation (b) and surface velocity field (c) in the catchment area of Ice Stream C. Also shown are the location of places referred to in the text and the location of the profile lines (gray lines) used in our basal energy model. The insert in (c) shows the approximate location of the graphs in Antarctica.

(BST) towards the now stopped trunk of Ice Stream C. The southern tributary originates from the southeastern divide in the Whitmore Mountains. Drawing its ice from the topographically higher Whitmore mountain block, the ice thickness of its sub-tributaries is thinner ($\sim 2000 \,\mathrm{m}$) and increases only when the ice is crossing a trench connecting a gap between the Whitmore mountain block and the Transantarctic Mountains with the emerging BST.

In our model, we consider two different flowlines for each tributary (Fig. 1; see Fig. 2 for details). For the northern tributary, one flowline (CN-NW) originates close to the junction of the Amundsen Sea embayment, the Weddell Sea and Ross Sea catchment areas. The other flowline (CN-S) initiates at the southeastern divide with the Weddell Sea catchment area, crossing the Byrd Plateau and flowing through a smaller sub-tributary into the main tributary. Both flowlines of the southern Ice Stream C Tributary (CS) originate in the Whitmore Mountains. The flow of the southern flowline (CS-S) recently got diverted towards Whillans Ice Stream. The northern flowline (CS-N) flows in its lower part, separated by a small ridge almost parallel to the southern flowline of the northern tributary. In the UpC area, flowline CN-NW passes the sticky spot in the north, while all other flowlines pass in the south.

THE BASAL ENERGY-BALANCE MODEL

We calculate the basal energy balance taking into account the geothermal flux (GT), the basal shear heating (SH), the heat (C) conducted through the ice away from the bed and the heat (M) used or released by either melting or freezing at the ice base:

$$M = SH + GT - C.$$
(1)

The melting or freezing rate (m) is the quotient of the energy balance and the latent heat of fusion (L) and density (ρ),

$$m = \frac{\mathrm{GT} + \tau_{\mathrm{b}} u_{\mathrm{b}} - K\beta_{\mathrm{b}}}{L\rho} \,. \tag{2}$$

The basal shear heating is the product of the basal shear stresses (τ_b) and the basal ice velocity (u_b) , while the heat conduction away from the bed is determined by the product of the basal temperature gradient (β_b) and the thermal conductivity of ice (K).

To obtain the basal temperature gradient, we use a numerical thermodynamical flowline model to model the temperature profile in the ice. The model includes heat conduction, vertical and horizontal advection and calculates the temperature profile in the ice. The temperature is calculated at 2l vertical nodes, repeated at 5 km horizontal intervals along the flowline, which is also the horizontal spacing for our basal energy-balance calculations.

As no data about basal conditions beneath ice-stream tributaries are available, we assume in the model that basal conditions in the ice-stream tributaries are similar to those in ice-stream trunks. Water is available for freeze-on at all times at the ice base. Basal temperatures are everywhere at the pressure-melting point (PMP; see Paterson, 1994, p. 212, for PMP calculation). Independent of the individual driving stress, basal shear stresses are taken to be low everywhere, between 1 and 10 kPa (Tulaczyk and others, 2000a; Kamb, 2001). In the model, we use a geothermal flux of 70 mW m⁻². H. Engelhardt (personal communication, 2002) calculated

this value using the temperature profile of Siple Dome. This is the only geothermal flux calculation in West Antarctica using a continuous temperature profile. Avoiding speculation about a possible distribution of the geothermal flux underneath the WAIS, we assume this value to be valid along all flowlines. The usage of a higher (lower) value would increase (decrease) basal melting and decrease (increase) basal freezing, but would not change the overall pattern of basal melting beneath the thick ice in the WAIS interior, and freezing in the shallower parts of the WAIS towards the ice streams.

Further on, we assume that basal sliding is the dominant part of ice movement (Engelhardt and Kamb, 1998; Kamb, 2001), allowing us to substitute the surface velocity (u_s) for the basal velocity in Equation (3) and to neglect internal deformation in the temperature model. Therefore we calculate temperature changes at a point within the ice sheet from the vertical heat conduction through the ice and the horizontal and vertical advection of heat (for more detailed discussion of Equation (5) and the effect of general assumptions, see, e.g., Robin, 1983; Paterson, 1994; Van der Veen, 1999).

$$\frac{\partial T}{\partial t} = k \frac{\partial^2 T}{\partial z^2} + u_{\rm h} \frac{\partial T}{\partial x} - u_{\rm v} \frac{\partial T}{\partial z}$$
(3)
$$\frac{dT}{dt} = k \frac{T_{z-1,x} - 2T_{z,x} + T_{z+1,x}}{dz^2} + u_{\rm h} \frac{T_{z,x} - T_{z,x-1}}{dx}$$

$$- A \frac{z}{H} \frac{T_{z-1,x} - T_{z+1,x}}{2dz},$$

where k is thermal diffusivity $(37.2 \text{ m}^2 \text{ a}^{-1} \text{ at } -1^\circ \text{C};$ Hooke, 1998), u_{h} is horizontal velocity (surface velocity), $u_{\text{v}} = AzH^{-1}$ is vertical velocity, A is accumulation rate, z is height over the bed, H is ice thickness, $T_{i,j}$ is temperature at depth *i* and location *j*, and *t* is time.

In several model runs, we have implemented a quadratic form of the parameterization of vertical ice velocity near the ice divide (Paterson 1994, p. 220) and have verified that the exact choice of parameterization has negligible effect on our calculations of basal freezing/melting.

We used an atmospheric lapse rate of 0.002 K m^{-1} to calculate the surface temperature distribution by assuming that they are linearly dependent on surface elevation. This value differs from the typical range of atmospheric lapse rates (0.006–0.010 K m⁻¹). It was calculated using modern 10 m temperatures at Byrd Station (-28°C at 1530 m a.s.l.; Robin, 1983) and Ice Stream C (-26°C at 500 m a.s.l.; Kamb, 2001; personal communication from H. Engelhardt, 2001). In analogy with surface temperature changes, we also parameterized spatial variation in accumulation rates along the simulated flowlines with surface elevation by using a linear fit to modern accumulation data (Giovinetto and others, 1990; Giovinetto and Zwally, 2000).

Surface and bed geometry were prescribed in our model using elevation and ice thickness from the BEDMAP dataset (Lythe and others, 2001), which was resampled to match the velocity dataset (Joughin and others, 1999). As the surface velocity field in the tributaries is close to the balance velocity, we assume the velocity field to be constant over time. Where Shabtaie and Bentley (1987) had mapped the buried shear margin of Ice Stream C, we gradually adjusted the velocity to 100 m a^{-1} as a transition to the faster ice-stream flow, assuming that the currently lower flow velocities are a result of the Ice Stream C stoppage and that flow in the past was similar to the balance velocity (Bamber and others, 2000). After initializing the model over 50 kyr with modern-day sur-



Fig. 3. Graphs showing the history of (a) surface temperature, (b) surface accumulation rate, and (c) WAIS elevation changes over the last 125000 years. Surface-temperature changes and parameterized accumulation rates are based on the Vostok ice-core record (Petit and others, 1999). Ice surface-elevation changes are based on our reinterpretation of the reconstructions by Jenssen (1983) and Steig and others (2001).

face temperatures, accumulation rates, ice velocities and icethickness values, we ran the model over a complete glacialinterglacial cycle from 125 kyr BP to today using a time-step of 10 model years. During the simulated glacial-interglacial cycle, we adjusted for temperature and accumulation changes (Fig. 3a and b) using the Vostok ice-core record (Petit and others, 1999). We smoothed the temperature and accumulation record using a 1000 year moving average. We also varied ice-surface elevation (Fig. 3c) using the WAIS elevation history estimated by Jenssen (1983) and Steig and others (2001). As the values given by Jenssen (1983) overestimate elevation changes, we adjust his numbers downward for the time period from 11500 years BP to 45000 years BP, using the difference between his value at 11500 years BP and the value for the same time given by Steig and others (2001). For simplicity, we assume uniform elevation changes along all flowlines and do not consider different thickness changes in different parts along the flowline (see, e.g., Steig and others, 200l; Waddington and others, http://igloo.gsfc.nasa.gov/



Fig. 4. Measured temperature profile at UpC camp (solid line; courtesy of H. Engelhardt); the temperature profiles of profile lines CN-S, CS-N and CS-S (dashed lines), originating on the Byrd Plateau and passing the UpC sticky spot in the south; and the temperature profile of profile line CN-NW (dotted line), which emerges from the BSB and passes the UpC sticky spot in the north. The basal temperature gradient of the three dashed lines is lower than the measured one, while the CN-NW profile has a higher temperature gradient.

Table 1. Basal temperature gradient (β), ice velocities (u) and basal freeze-on rates (m) at 1kPa

Location	eta	u	m
	${ m K}~{ m km}^{-1}$	${ m m~a}^{-1}$	$\mathrm{mm}\mathrm{a}^{-1}$
CN-NW	54.6	8	4.5
CN-S CS-N CS-S	48.6 47.3 46.6	12	3.2 2.9 2.8
Sticky spot	52.0	2	4.0
Ice Stream C UpC 2002	53.0 54.5	12	4.2 4.5

wais/abstracts0l/Waddington.html). If we were to do so, we would have to engage in speculation since there are no data that could be used to infer ice-thickness histories in different parts of our flowlines. Also several model runs have showed that possibly larger ice-thickness changes in the UpC area or ice-stream trunks (Waddington and others, http://igloo.gsfc. nasa.gov/wais/abstracts0l/Waddington.html) have negligible influence on the basal ice production upstream of this location. This also can be easily illustrated, as the ice influenced by a thickened Last Glacial Maximum (LGM) UpC area has moved further downstream and the ice currently present in the UpC area has moved into the UpC area from further upstream.

RESULTS

Figure 4 compares the modeled present-day temperature profile at the UpC sticky spot (around km 600 in the profile lines) with the measured one (personal communication from H. Engelhardt, 2001). Along the flowline, the simulated basal



Fig. 5. Results for the profile lines CN-NW (a-f) and CN-S (g-l) of the northern Ice Stream C tributaries. (a, g) Current (solid line) and early-Holocene (10 kyr; dashed line) basal temperature gradient; (b, c, h, i) basal melt and freezing cumulative along the flowline starting at the ice divide, using 1kPa (solid line) and 10 kPa (dashed line) as basal shear stress; (d, j) basal freezing rates along flowline using 1kPa (solid line) and 10 kPa (dashed line) as basal shear stress; (e, k) basal freezing rates for present day (1kPa; solid line) and early Holocene (10 kyr; 1kPa dashed line, 10 kPa dash-dotted line); (f, l) time evolved since ice started to move away from ice divide.



Fig. 6. Results for the profile lines CS-N (a-f) and CS-S (g-l) of the northern Ice Stream C tributaries. (a, g) current (solid line) and early-Holocene (10 kyr; dashed line) basal temperature gradient; (b, c, h, i) basal melt and freezing cumulative along the flowline starting at the ice divide, using 1 kPa (solid line) and 10 kPa (dashed line) as basal shear stress; (d, j) basal freezing rates along flowline using 1 kPa (solid line) and 10 kPa (dashed line) as basal shear stress; (e, k) basal freezing rates for present day (1 kPa; solid line) and early Holocene (10 kyr; 1 kPa dashed line, 10 kPa dash-dotted line); (f, l) time evolved since ice started to move away from ice divide.

temperature gradient ranges from 30 to 60 K km^{-1} (Figs 5a and g and 6a and g). The flowlines passing south of the sticky spot show a similar basal temperature gradient of 47.5 K km⁻¹

 $(\pm l\,K\,km^{-1})$ (see Table 1). This is 6.2 K km^{-1} lower than the measured temperature gradients in the southern Ice Stream C branch. The basal temperature gradient for the northern

Table 2. Modeled basal melting and freezing (m) beneath Ice Stream C tributaries

	Basal shear stress				
Flowline	1 kPa		10 kPa		
	Melting	Freezing	Melting	Freezing	
CN-NW	33	17	44	7	
CN-S	15	13	26	4	
CS-N	1	17	9	4	
CS-S	9	11	19	4	

Ice Stream C branch at the sticky spot (profile line CN-NW) is 54.6 K km⁻¹: similar in magnitude to, yet 7.1 K km⁻¹ higher than, those for the southern Ice Stream C branch. Unfortunately, no measured temperature profile is available for the northern Ice Stream C branch, so we focus our comparison of basal ice accumulation below on the profile lines passing through the southern ice-stream branch, in which the discovery was made. In the southern ice-stream branch at the UpC location, the estimated modern freeze-on rate is $4.2 \pm 0.2 \text{ mm a}^{-1}$ (Table 1; Kamb, 2001), ~1.3 mm a⁻¹ higher than the modeled rates. Hence the model underestimates basal freeze-on rates at this location by ~30%.

Along the flowlines, the threshold basal temperature gradient for basal freezing varies, dependent on ice velocity, from 33 to 35 K km⁻¹ (33 to 45 K km⁻¹ for $\tau_{\rm b} = 10$ kPa), for a basal shear stress of 1 kPa. Comparing the obtained melting rates with the freezing rates shows a distinct difference between the northern and southern tributaries (Table 2; Figs 5b-e and h-k and 6b-e and h-k). Basal melting is high underneath the thick ice of the WAIS interior (over the first 200-300 km of the flowlines), while basal freezing is dominant in the thinner parts of the tributaries closer to the UpC area. Basal melting is particularly high in the BSB (km 50-200, flowline CN-NW, Fig. 5a-f) and starts to decrease where ice emerges from the BSB. Further thinning of ice in the BST leads to a steepening of the basal temperature gradient and to basal freezing, allowing the accretion of up to 17 m of basal ice in the model.

Basal freezing close to the ice divide in the tributaries originating on the Byrd Plateau (southern tributaries and CN-S; see Fig. l) possibly indicates that the ice sheet is frozen to its bed and that heat used in the model to melt ice at the base of the ice sheet further downstream will actually be used to warm the basal ice until the PMP is reached.

Taking into account the possibility of a frozen bed on the Byrd Plateau, we consider only the last 300–400 km of each flowline while estimating the thickness of basal ice at the UpC sticky spot (100 km from the end of each flowline). Considering the underestimation of basal freeze-on rates (~1.3 mm a⁻¹ or ~30 %) in the southern ice-stream branch at the UpC sticky spot, the modeled basal ice thickness (6– 13 m) at the UpC sticky spot (Table 3) is similar in magnitude to the observed basal ice thickness of 10–25 m. With a higher basal shear stress of 10 kPa, the model still obtains a basal ice thickness of up to 3 m at this location; while with a basal shear stress of 22 kPa no basal ice is obtained. A further increase in basal shear stress to 70 kPa leads to > 150 m of basal melt.

These calculations illustrate the sensitivity of predicted basal ice thickness at UpC to the assumed magnitude of

Table 3. Modeled basal ice thickness (m) at the UpC sticky spot (km 600)

Flowline	Basal sh	ear stress
	1 kPa	10 kPa
CN-NW	10	3
CN-S	6	1
CS-N	13	3
CS-S	6	2

basal shear stress. In reality, there is likely to be a spatial variability in the basal shear stress, but, with no firm observational constraints available, we opted for a simple assumption of a constant basal shear stress along all of the considered flowlines. In the case of advection of the basal ice layer from upstream (see third paragraph of Discussion below), these sensitivity analyses also show that moderate basal shear stresses (~ 20 kPa) have to be present in the ice-stream tributaries in order to preserve such a feature.

DISCUSSION

If our conceptual picture of basal ice formation by basal freeze-on upstream of UpC is correct, we have to contend with the fact that the model predicts that basal melting and freezing occur in spatially separated areas. In order to sustain the basal freeze-on expected from our model results, either water is transported between these two areas, or basal water, locally stored in the system, from earlier melt events is used in the freeze-on process, or both. Although we do not know precisely what form the latter water storage of subglacial water would take (subglacial lakes?), one can observe that the basal temperature gradients in our model (Figs 5a and g and 6a and g; solid line present day, dashed line 10 kyr BP) are consistently smaller during the late LGM and early Holocene than today. This model prediction suggests the possibility of more widespread melting and higher melting rates during this time. The ultimate physical reason for the lower late-LGM and early-Holocene basal temperature gradient is the assumed greater ice thickness at that time (Steig and others, 2001). We speculate that this higher basal water production may have also helped to lubricate the base of the ice sheet and promote its rapid Holocene decay (Conway and others, 1999).

One can use our results to consider a regional basal water balance over roughly the area of the BSB and other similarly deep parts of the WAIS interior (Parizek and others, 2002) as well as the shallower part of the Ice Stream C tributaries. This regional balance between basal melting and basal freezing for the Ice Stream C tributaries (see Table 2) may have shifted from being a source of basal water in the early Holocene to a sink in modern times. Reduction in the amount of water stored in the subglacial hydrological system could increase basal resistance (Tulaczyk and others, 2000b), which would be consistent with present-day force-budget calculations (Joughin and others, 2002) and the observed slow-down of Whillans Ice Stream (Joughin and others, 2002) as well as the shut-down of Ice Stream C (Anandakrishnan and Alley, 1997). It is important to note, however, that the conjectured persistence of high basal shear stresses (~70 kPa; Joughin and others, 2002) in the past along the entire flowline would

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have caused extensive basal melting (150 m in our model), making basal freeze-on an unlikely explanation for the observed basal ice layer at the UpC sticky spot.

The model currently does not take into account the possibility of a frozen bed in parts of the WAIS interior and also does not treat explicitly latent-heat transport through basal water drainage. Further uncertainties of the model result from unknown changes in the ice-flow history and in the WAIS elevation history. Given the simplifying assumptions that went into the construction of our model, we cannot use the model results to claim that the debris-laden basal ice layer found in boreholes drilled in the UpC area of Ice Stream C must have formed by freeze-on beneath the tributaries of this ice stream. We can only say that such genesis of the basal layer is admissible, given a restrictive condition of very low basal shear stresses everywhere beneath simulated flowlines. An alternative explanation for the origin of this layer is suggested by Christoffersen and Tulaczyk (2003), who show that $\sim 20 \,\mathrm{m}$ of frozen-on basal ice may have formed beneath interstream ridges as a result of the cooling trend during the last glacial cycle. Kamb's (2001) analysis of borehole data from interstream ridges indicates that they may be underlain by thick (up to 10-41m) debris-laden basal ice layers. An ice core recovered from Byrd Station contained ~ 5 m of debris-laden basal ice. Hence, an alternative explanation for the basal layer at UpC is that it represents a feature formed in slow-moving ice (outside of icestream tributaries), which was then incorporated into the tributaries of Ice Stream C and advected to the UpC area.

Future field research, in particular the drilling of boreholes to examine ice-stream tributary beds in the WAIS interior and additional temperature measurements, seems necessary for constraining basal water production rates and basal water distribution beneath ice-stream tributaries and the WAIS interior. Improved, physically based parameterizations of these processes are needed to accurately model the past and future behavior of the WAIS (Steig and others, 2001) and to assess the possible threat of a rapid sealevel rise caused by a rapid disintegration of the WAIS (Mercer, 1978; Scherer and others, 1998).

CONCLUSIONS

Our numerical model applied to three different flowlines converging near the UpC sticky spot shows that the 10-25 m thick basal ice layer observed in the UpC area may have formed by basal freeze-on occurring beneath ice-stream tributaries. Admittedly, the model can produce a sufficient amount of basal ice only if we assume that beds of tributaries are as weak as the till beds found beneath ice-stream trunks (1-10 kPa). In our model, basal melting was highest during the LGM, due to a thickened ice sheet. According to the model, basal melting rates in the WAIS interior peaked around 15 kyr BP. This increase in availability of basal water could have played a role in initiating the rapid Holocene decay of the WAIS. A subsequent increase in basal freeze-on rates related to ice thinning might have shifted the basal energy balance towards freezing. The latter process may have gradually reduced the amount of basal water in the subglacial system over the course of the Holocene. A reduction of basal water availability can lead to an increase in basal resistance to fast ice motion and should be considered as one possibility for the recent shut-down of Ice Stream C and the current

slow-down of Whillans Ice Stream. These conjectures, however, need to be verified by field observations.

ACKNOWLEDGEMENTS

We thank all members of the research team participating in the Sticky Spot field campaign 2000/01, in particular B. Kamb and H. Engelhardt for leading this field campaign and for giving access to samples and data of this and former field campaigns. The basal ice layer at UpC was imaged using a unique borehole camera designed and prepared by the Caltech Jet Propulsion Laboratory team, which includes F. Carsey, H. Engelhardt, A. Behar, R. Bolsey, K. Boykins, R. Ivlev, L. Lane and K. Manatt. We are grateful for the thorough review from D. Reusch, J. Johnson and E. Waddington and the helpful comments of H. Engelhardt. We thank also the U.S. National Science Foundation (NSF) Office of Polar Programs (OPP) for logistical support of this and previous field campaigns. This material is based upon work supported by the NSF under grants NSF-OPP 9873593 and NSF-OPP 0096302.

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