A UNIMODAL PROPERTY OF PURELY IMAGINARY ZEROS OF BESSEL AND RELATED FUNCTIONS

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ABSTRACT. We show, among other things, that, for n = 0, 1, the negative of the square of a purely imaginary zero of $J_{\nu}^{(n)}(x)$ is unimodal on (n - 2, n - 1). One of the important tools in the proof is the Mittag-Leffler partial fractions expansion of $J_{\nu+1}(z)/J_{\nu}(z)$.

1. Introduction. For n = 0, 1, 2 the smallest positive zero of $J_{\nu}^{(n)}(x)$ decreases to 0 as ν decreases to n - 1. For n = 0, 1, this is a classical result [14, Chapter 15], while for n = 2 it is very recent; see [10], [12], [13], [15]. As ν decreases below n - 1 the zero becomes purely imaginary first moving away from the origin and then returning to the origin (along the imaginary axis) as ν approaches n - 2. This behaviour was observed numerically by Kerimov and Skorokhodov ([8] and [9]). It can be described by saying that the negative of the square of such a purely imaginary zero is *unimodal* on (n - 2, n - 1). Curiously, this was proved analytically [4] first for the case n = 2, *i.e.*, for a purely imaginary zero of $J_{\nu}''(x)$ on (0, 1). Our main purpose here is to deal with the cases n = 0, 1, i.e., to prove the corresponding properties of the purely imaginary zeros of $J_{\nu}(x)$ and $J_{\nu}'(x)$ on (-2, -1) and (-1, 0) respectively. We also give a slightly simpler version of the proof in the case n = 2. The case n = 0 is illustrated in Figure 1.

We will have need of the Bessel differential equation

(1.1)
$$z^2 J_{\nu}''(z) + z J_{\nu}'(z) + (z^2 - \nu^2) J_{\nu}(z) = 0,$$

the recurrence relations [14, p. 45]

(1.2)
$$J_{\nu-1}(z) + J_{\nu+1}(z) = \frac{2\nu}{z} J_{\nu}(z),$$

and

(1.3)
$$zJ'_{\nu}(z) - \nu J_{\nu}(z) = -zJ_{\nu+1}(z),$$

as well as the Mittag-Leffler type expansion [14, p. 498]

(1.4)
$$\frac{J_{\nu+1}(z)}{J_{\nu}(z)} = \sum_{n=1}^{\infty} \frac{2z}{j_{\nu n}^2 - z^2},$$

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where $\{\pm j_{\nu n}\}$ is the sequence of zeros of the entire function $z^{-\nu}J_{\nu}(z)$. We note that these zeros are all real if $\nu > -1$, and according to the conventional notation [14, p. 497], $0 < j_{\nu 1} < j_{\nu 2} < \cdots$.

We use the method of the last part of [4], *i.e.*, we prove the unimodality of a function by showing that it is concave down at every turning point, hence there can be only one turning point.

2. A general function. All of the functions which we deal with can be subsumed in the general formula

$$f_{\nu}(z) = \left(cz^2 + f(\nu)\right)J_{\nu}(z) - azJ_{\nu+1}(z) + bzJ_{\nu-1}(z).$$

The most important special cases are:

- (i) $f_{\nu}(z) = zJ_{\nu-1}(z)$, got by taking b = 1, $a = c = f(\nu) = 0$;
- (ii) $f_{\nu}(z) = z J'_{\nu}(z)$, got by taking $a = 1, b = c = 0, f(\nu) = \nu$ and using (1.3);
- (iii) $f_{\nu}(z) = \alpha J_{\nu}(z) + z J'_{\nu}(z)$, got by taking $a = 1, b = c = 0, f(\nu) = \alpha + \nu$;
- (iv) $f_{\nu}(z) = -z^2 J_{\nu}''(z)$, got by taking $a = c = 1, b = 0, f(\nu) = \nu \nu^2$ and using (1.1) and (1.3).

THEOREM 2.1. Let $\nu > -1$ and let $i\rho$ be a purely imaginary zero of $f_{\nu}(z)$. Then

(2.1)
$$\lambda(\nu)\frac{d\rho^2}{d\nu} = \mu(\nu),$$

where

(2.2)
$$\lambda(\nu) = -\frac{f(\nu) + 2b\nu}{\rho^2} - 2(b+a)\sum_{n=1}^{\infty} \frac{\rho^2}{(j_{\nu,n}^2 + \rho^2)^2},$$

or

(2.3)
$$\lambda(\nu) = -c + 2(b+a) \sum_{n=1}^{\infty} \frac{j_{\nu,n}^2}{(j_{\nu,n}^2 + \rho^2)^2},$$

and

(2.4)
$$\mu(\nu) = 4(b+a)\rho^2 \sum_{n=1}^{\infty} \frac{j_{\nu,n} dj_{\nu n} / d\nu}{(j_{\nu,n}^2 + \rho^2)^2} - f'(\nu) - 2b.$$

For those values of ν for which $d\rho^2/d\nu = 0$, we have

(2.5)
$$\lambda(\nu)\frac{d^2\rho^2}{d\nu^2} = \rho^2\mu_1(\nu),$$

where

(2.6)

$$\mu_1(\nu) = 4(b+a) \left[\sum_{n=1}^{\infty} \frac{j_{\nu,n} d^2 j_{\nu n} / d\nu^2}{(j_{\nu,n}^2 + \rho^2)^2} + \sum_{n=1}^{\infty} \frac{(dj_{\nu n} / d\nu)^2}{(j_{\nu,n}^2 + \rho^2)^2} - 4 \sum_{n=1}^{\infty} \frac{(j_{\nu n} dj_{\nu n} / d\nu)^2}{(j_{\nu,n}^2 + \rho^2)^3} \right] - f''(\nu) / \rho^2.$$

PROOF. The equation

$$f_{\nu}(z) = \left(cz^2 + f(\nu)\right)J_{\nu}(z) - azJ_{\nu+1}(z) + bzJ_{\nu-1}(z) = 0$$

can be put in the form

$$cz^{2} + f(\nu) - az \frac{J_{\nu+1}(z)}{J_{\nu}(z)} + bz \frac{J_{\nu-1}(z)}{J_{\nu}(z)} = 0,$$

except at the zeros of $J_{\nu}(z)$.

Using (1.2) and (1.4) this becomes

$$cz^{2} + f(\nu) - 2(b+a) \sum_{k=1}^{\infty} \frac{z^{2}}{j_{\nu k}^{2} - z^{2}} + 2b\nu = 0.$$

For $z = i\rho$ we get

(2.7)
$$c - \frac{f(\nu) + 2b\nu}{\rho^2} = 2(b+a)\sum_{k=1}^{\infty} \frac{1}{j_{\nu k}^2 + \rho^2}.$$

Differentiating this equation with respect to ν and multiplying by ρ^2 , we get (2.1), with $\lambda(\nu)$ given by (2.2) and $\mu(\nu)$ by (2.4). Then, using (2.7), we can express $\lambda(\nu)$ in the form (2.3). Differentiating (2.1) and using $d\rho^2/d\nu = 0$, we get (2.5), where $\mu_1(\nu)$ is given by (2.6).

In order to justify the term-by-term differentiation we have to verify that the differentiated series or, equivalently, all the infinite series in (2.3), (2.4) and (2.6), converge uniformly in ν , in any closed subinterval $[\nu_0, \nu_1]$ of $(-1, \infty)$. In the case of the series in (2.3), (2.4) and the second and third series in (2.6), this is a consequence of the inequality [11, p. 471]

$$(2.8) \qquad (\nu+1)\frac{dj_{\nu k}}{d\nu} \le j_{\nu k}$$

and the convergence of

$$\sum_{k=1}^{\infty} j_{\nu k}^{-2}.$$

To deal with the first infinite series in (2.6), we use the representation [1, p. 87]

(2.9)
$$j'' = 2 \int_0^\infty K_0(2j \sinh t) e^{-2\nu t} I(\nu, t) dt$$

where

(2.10)
$$I(\nu, t) = 2\nu j' \tanh t + j' \tanh^2 t - 2jt.$$

Here we are using the notation $j = j_{\nu k}$ and the primes denote differentiation with respect to ν . For the uniform convergence of the first infinite series in (2.6), it is sufficient to show that

(2.11)
$$\frac{d^2 j_{\nu k}}{d\nu^2} < F(\nu) j_{\nu k}, \quad \nu_0 \le \nu \le \nu_1,$$

where $F(\nu)$ is bounded on $\nu_0 \le \nu \le \nu_1$. Now, from (2.10), j'' is bounded by

(2.12)
$$j'(2|\nu|+1)2\int_0^\infty K_0(2j\sinh t)e^{-2\nu t}\,dt+4j\int_0^\infty K_0(2j\sinh t)e^{-2\nu t}\,dt.$$

Using [14, p. 508], (2.13)

and (2.8), we find that the first term here is bounded by $(2|\nu| + 1)(\nu + 1)^{-2}j$. Using $j_{\nu k} > \nu + 1$, $\sinh t > t$ and the decreasing character of $K_0(t)$ as a function of t, t > 0, we find that the second term is bounded by

 $j' = 2j \int_0^\infty K_0(2j\sinh t)e^{-2\nu t} dt,$

$$4j \int_0^\infty K_0 \Big(2(\nu+1)t \Big) e^{-2\nu t} t \, dt.$$

Thus (2.11) holds and this completes the verification of the validity of the differentiation and the proof of Theorem 2.1.

3. Zeros of $J_{\nu}(z)$. We first deal with the case $f_{\nu}(z) = zJ_{\nu-1}(z)$, got by taking b = 1, $a = c = f(\nu) = 0$. It is well known [14, p. 483] that, for $-2 < \nu < -1$, $J_{\nu}(z)$ has a pair of purely imaginary zeros.

THEOREM 3.1. Let $\pm i\rho$ be the purely imaginary zeros of $J_{\nu}(z)$ for $-2 < \nu < -1$. Then ρ^2 is unimodal on (-2, -1).

PROOF. We will actually work with the zeros of $J_{\nu-1}(z)$ on (-1, 0). From §2, we have (2.1) with

(3.1)
$$\lambda(\nu) = 2 \sum_{n=1}^{\infty} \frac{j_{\nu,n}^2}{(j_{\nu,n}^2 + \rho^2)^2},$$

and

(3.2)
$$\mu(\nu) = 4\rho^2 \sum_{n=1}^{\infty} \frac{j_{\nu,n} dj_{\nu n} / d\nu}{(j_{\nu,n}^2 + \rho^2)^2} - 2.$$

CASE (i): $-1 < \nu \le -0.8$. In this range we will show that $d\rho^2/d\nu > 0$, so that ρ^2 is increasing. Since $\lambda(\nu) > 0$, we need to show that $\mu(\nu) > 0$ for this range of values of ν . We will need the following results [5]:

(3.3)
$$j_{\nu n} \frac{dj_{\nu n}}{d\nu} > 4 - \frac{8(\nu+1)(\nu+3)}{j_{\nu n}^2} + \frac{32(\nu+1)^2(\nu+2)^2}{j_{\nu n}^4}, \quad \nu > -1;$$

(3.4)
$$4(\nu+1) < j_{\nu_1}^2 < 4(\nu+1)(\nu+2), \quad \nu > -1.$$

We also need inequalities for $j_{\nu 1}^2$ in the case $-2 < \nu < -1$, when it is negative. Some simple bounds in this case are [7]

(3.5)
$$2(\nu+1)(\nu+3) < j_{\nu 1}^2 < 4(\nu+1)(\nu+2), \quad -2 < \nu < -1.$$

or, in terms of our present notation,

(3.6)
$$-4\nu(\nu+1) < \rho^2 < -2\nu(\nu+2), \quad -1 < \nu < 0.$$

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We have

(3.7)
$$\mu(\nu) = -2 + 4\rho^2 \sum_{n=1}^{\infty} \frac{j_{\nu,n} dj_{\nu n} / d\nu}{j_{\nu,n}^4} \left[1 + \frac{\rho^2}{j_{\nu n}^2} \right]^{-2}.$$

Clearly

$$\left[1+\frac{\rho^2}{j_{\nu n}^2}\right]^{-1} > \left[1+\frac{\rho^2}{j_{\nu 1}^2}\right]^{-1} > 1-\frac{\rho^2}{j_{\nu 1}^2} > 1+\frac{2\nu(\nu+2)}{j_{\nu 1}^2} > 1+\frac{\nu(\nu+2)}{2(\nu+1)},$$

where we have used the upper bound in (3.6) and the lower bound in (3.4). On the other hand, from (3.3),

$$\sum_{n=1}^{\infty} \frac{j_{\nu n} dj_{\nu n} / d\nu}{j_{\nu n}^4} > 4\sigma_{\nu}^{(2)} - 8(\nu+1)(\nu+3)\sigma_{\nu}^{(3)} + 32(\nu+1)^2(\nu+2)^2\sigma_{\nu}^{(4)}.$$

Here we have used the notation

$$\sigma_{\nu}^{(n)} = \sum_{k=1}^{\infty} j_{\nu k}^{-2n}$$

and the closed form expressions for these sums in [14, p. 502]. This gives

$$\sum_{n=1}^{\infty} \frac{j_{\nu n} dj_{\nu n} / d\nu}{j_{\nu n}^4} > \frac{5\nu + 11}{8(\nu + 1)^2(\nu + 3)(\nu + 4)}.$$

Using the above bounds and the lower bound for ρ^2 in (3.6), we get from (3.7)

$$\mu(\nu) > -2 - \frac{\nu(\nu^2 + 4\nu + 2)^2(5\nu + 11)}{2(\nu + 1)^3(\nu + 3)(\nu + 4)},$$

and it is easy to check that this is positive for $-1 < \nu \leq -0.8$.

CASE (ii): $-0.8 < \nu < 0$. Here we show that $d^2 \rho^2 / d\nu^2$ is negative at the turning points of ρ^2 . In the present case (2.6) can be written

(3.8)
$$\mu_1(\nu) = 4 \left[\sum_{n=1}^{\infty} \frac{j_{\nu n} d^2 j_{\nu n} / d\nu^2}{(j_{\nu n}^2 + \rho^2)^2} + \sum_{n=1}^{\infty} \frac{(d j_{\nu n} / d\nu)^2}{(j_{\nu n}^2 + \rho^2)^2} - 4 \sum_{n=1}^{\infty} \frac{(j_{\nu n} d j_{\nu n} / d\nu)^2}{(j_{\nu n}^2 + \rho^2)^3} \right].$$

Now the first term on the right here is negative [1] and the sum of the two remaining terms will certainly be negative provided that

(3.9)
$$\rho^2 - 3j_{\nu 1}^2 < 0.$$

But

$$\rho^2 - 3j_{\nu 1}^2 < -2(\nu^2 + 8\nu + 6)$$

from (3.4) and (3.6) and this is certainly negative in case $-0.8 < \nu < 0$. Thus the second derivative of ρ^2 with respect to ν is negative at points where the first derivative is 0; hence there can be only one such point and it is a relative maximum. This proves the unimodal property and completes the proof of Theorem 3.1.



FIGURE 1: $j_{\nu 1}^2$ vs. ν

In Figure 1, we give the graph of $j_{\nu 1}^2$ versus ν for $-2 \le \nu \le 0$. This graph, based on a computation described in [7], suggests strongly that $j_{\nu 1}^2$ is convex on $-2 < \nu < 0$. It is shown [2] that it is convex on $(0, \infty)$ and conjectured that the convexity extends to $(-1, \infty)$.

A numerical calculation indicates that the smallest value of $j_{\nu 1}^2$ is -1.6075 to 5 digit accuracy and it occurs for ν between -1.698 and -1.697.

4. Dini functions and derivatives of Bessel functions. Here we deal with the case $f_{\nu}(z) = H_{\nu}(z) = \alpha J_{\nu}(z) + z J'_{\nu}(z)$, got by taking a = 1, b = c = 0, $f(\nu) = \alpha + \nu$. We call these Dini functions because they arise in expansions due to Dini [14, Chapter 18]. In the special case $\alpha = 0$, we are, of course, dealing with the zeros of $J'_{\nu}(z)$. There do not appear to be many results in the literature on the monotonicity of purely imaginary zeros $\pm i\rho$ of $H_{\nu}(z)$, though it is shown in [6, pp. 78–79] that if $\alpha < 0$, then ρ^2 is decreasing on $(0, -\alpha)$. We will prove:

THEOREM 4.1. Let $H_{\nu}(z, \alpha) = \alpha J_{\nu}(z) + z J'_{\nu}(z)$, where $-1/2 \leq \alpha < 1$ and $-1 < \nu < -\alpha$. $H_{\nu}(z, \alpha)$ has a pair of purely imaginary zeros $\pm i\rho(\nu, \alpha)$. $\rho^2(\nu, \alpha)$ is unimodal on $(-1, -\alpha)$, i.e., there exists a number $\nu_0(\alpha)$ such that $\rho^2(\nu, \alpha)$ increases on $(-1, \nu_0(\alpha))$ and decreases on $(\nu_0(\alpha), -\alpha)$.

COROLLARY. If $\pm i\rho$ are purely imaginary zeros of $J'_{\nu}(z)$ then ρ^2 is unimodal on (-1,0).

PROOF OF THEOREM 4.1. The question of existence of such zeros is equivalent to the question of whether equation (2.7) which is, in this case,

(4.1)
$$-(\nu+\alpha) = 2\sum_{n=1}^{\infty} \frac{1}{j_{\nu n}^2/\rho^2 + 1},$$

can be satisfied. The right-hand side here increases from 0 to ∞ as ρ increases from 0 to ∞ , whereas the left-hand side remains constant and positive. Thus the existence of $\rho(\nu, \alpha)$ is clear. It is also clear from

(4.2)
$$\rho^2 < -\frac{\alpha + \nu}{2 + \alpha + \nu} j_{\nu 1}^2$$

[6, (3.2)] that $\rho(\nu, \alpha)$ vanishes as $\nu \to -\alpha^-$ and, since $j_{\nu 1} \to 0$, it also vanishes as $\nu \to -1^+$. In the present case, (2.5) holds with $\lambda(\nu)$ and $\mu_1(\nu)$ given by (3.1) and (3.8), but with ρ interpreted as in the current section. Now, as in §3, the first term on the right of (3.8) is negative [1] and the sum of the two remaining terms will certainly be negative provided that

(4.3)
$$\rho^2 - 3j_{\nu_1}^2 < 0.$$

But this follows from (4.2). Thus the second derivative of ρ^2 with respect to ν is negative at points where the first derivative is 0; hence there can be only one such point and it is a relative maximum. This proves the unimodal property.

Unfortunately, it does not seem to be possible to handle the case where $\alpha < -1/2$ in this way. The problem is that the inequality (4.3) seems to break down in this case. This is because $j_{\nu 1}^2 \sim 4(\nu + 1)$ and $\rho^2 \sim -4(\nu + 1)(\alpha - 1)/(\alpha + 1)$ as $\nu \rightarrow -1^+$ [7].

With regard to the Corollary, it is of interest to point out that if $j'_{\nu 1}$ denotes the purely imaginary zero of $J'_{\nu}(x)$, then the smallest value of $j'^2_{\nu 1}$ is -0.60602 to 5 digit accuracy and it occurs for ν between -0.5699 and -0.5696.

REMARK. Since the purely imaginary zeros of $J_{\nu}(z)$ are real zeros of $I_{\nu}(z)$, we may restate the results on $j'_{\nu 1}$ in the following way: The unique positive zero of $I'_{\nu}(x)$ increases from 0 to i_0 (= 0.7759 to four digits) as ν increases from -1 to $\nu_0(0)$ ($-0.5699 < \nu_0(0) <$ -0.5697) and then decreases again to 0 as ν increases from $\nu_0(0)$ to 0.

5. Zeros of $J''_{\nu}(z)$. Here we discuss the case $f_{\nu}(z) = -z^2 J''_{\nu}(z)$, got by taking a = c = 1, b = 0 and $f(\nu) = \nu - \nu^2$. In this situation, we have (2.1) and (2.5) where it is better to leave $\lambda(\nu)$ in the form given by (2.2):

(5.1)
$$\lambda(\nu) = \frac{\nu^2 - \nu}{\rho^2} - 2\sum_{n=1}^{\infty} \frac{\rho^2}{(j_{\nu n}^2 + \rho^2)^2},$$

and the function $\mu_1(\nu)$ in (2.5) is given by

(5.2)
$$\mu_1(\nu) = 4 \left[\sum_{n=1}^{\infty} \frac{j_{\nu n} d^2 j_{\nu n} / d\nu^2}{(j_{\nu n}^2 + \rho^2)^2} + \sum_{n=1}^{\infty} \frac{(dj_{\nu n} / d\nu)^2}{(j_{\nu n}^2 + \rho^2)^2} - 4 \sum_{n=1}^{\infty} \frac{(j_{\nu n} dj_{\nu n} / d\nu)^2}{(j_{\nu n}^2 + \rho^2)^3} \right] + \frac{2}{\rho^2}.$$

We note that $\lambda(\nu) < 0, 0 < \nu < 1$, so in order to get the unimodality of $j_{\nu 1}^{\prime\prime 2}$, we need to show that $\mu_1(\nu) > 0$ in this interval. Apart from a change in notation (ν being replaced by $-\nu$) the approach and formulas here agree with those in [4]. However, we can simplify the proof given in [4] by noting that the first two terms in (5.2), when combined, are equal to

$$2\sum_{n=1}^{\infty} \frac{(d^2 j_{\nu n}^2 / d\nu^2)}{(j_{\nu n}^2 + \rho^2)^2}$$

which is clearly positive since $d^2 j_{\nu n}^2 / d\nu^2 > 0$, $\nu > 0$ [2]. Thus it remains to show that the sum of the remaining two terms is positive, *i.e.*, that

(5.3)
$$8\sum_{n=1}^{\infty} \frac{(j_{\nu n} dj_{\nu n} / d\nu)^2}{(j_{\nu n}^2 + \rho^2)^3} < \frac{1}{\rho^2}.$$

But from [3, (1.5)], we have

(5.4)
$$\frac{1}{\rho^2} > \frac{2\nu + 1}{2\nu(1 - \nu^2)}.$$

while the inequality (2.8) and the Rayleigh sum [14, p. 502]

$$\sum_{n=1}^{\infty} j_{\nu n}^{-2} = 1 / [4(\nu + 1)]$$

show that the left-hand side of (5.3) is $< 2/(\nu + 1)^3$. It is a simple matter to show that

$$\frac{2}{(\nu+1)^3} < \frac{2\nu+1}{2\nu(1-\nu^2)}$$

for $0 < \nu < 1$, so we find that $\mu_1(\nu) > 0$, $0 < \nu < 1$ and this completes the proof.

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